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# Spectral Theory Sum Rules, Meromorphic Herglotz Functions and Large Deviations

Barry Simon

IBM Professor of Mathematics and Theoretical Physics, Emeritus  
California Institute of Technology  
Pasadena, CA, U.S.A.



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My goal in this talk is to explain two sum rules in the spectral theory of orthogonal polynomials, one of which I was involved with about 15 years ago.

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# Szegő Recursion and Verblunsky Coefficients

We start with orthogonal polynomials on the unit circle, aka OPUC. Let  $d\mu$  be a probability measure on  $\partial\mathbb{D}$ . Then, there are monic orthogonal polynomials,  $\Phi_n$ , and recursion relations due to Szegő in 1939

$$\Phi_{n+1}(z) = z\Phi_n(z) - \bar{\alpha}_n\Phi_n^*(z)$$

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where  $\Phi_n^*$  is the polynomial obtained by conjugating and reversing the order of the coefficients. The  $\{\alpha_n\}_{n=0}^\infty$ , called *Verblunsky coefficients*, lie in  $\mathbb{D}$  and there is a one-one correspondence, called the *Verblunsky map*, from measures of infinite support and sequences in  $\mathbb{D}$ .

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# Szegő's Theorem: Toeplitz version

Szegő's Theorem concerns probability measures on  $\partial\mathbb{D}$  of the form

$$d\mu(\theta) = w(\theta) \frac{d\theta}{2\pi} + d\mu_s(\theta)$$

where  $d\mu_s$  is singular w.r.t.  $d\theta$ .

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$$c_{k\ell} \equiv \int e^{i(k-\ell)\theta} d\mu(\theta) = \langle e^{-ik\cdot}, e^{-i\ell\cdot} \rangle_{L^2(d\mu)}$$

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In 1915, Szegő proved that

$$\lim_{n \rightarrow \infty} D_n(d\mu)^{1/n} = \exp \left[ \int \log(w(\theta)) \frac{d\theta}{2\pi} \right]$$

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While this is true in general, Szegő only proved it when  $d\mu_s = 0$ .

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# Szegő's Theorem: OPUC version

In 1920, Szegő realized that, because a Toeplitz matrix is just the Gram matrix of  $\{z^j\}_{j=0}^{n-1}$ , it is also the Gram matrix of  $\{\Phi_j\}_{j=0}^{n-1}$  which is diagonal so

$$D_n = \prod_{j=0}^{n-1} \|\Phi_j\|^2$$

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so using that  $\|\Phi_j\|$  is monotone decreasing (by a variational argument), one has an equivalent form of his theorem, namely

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$$\lim_{n \rightarrow \infty} \|\Phi_n\|^2 = \exp \left[ \int \log(w(\theta)) \frac{d\theta}{2\pi} \right]$$

But the recursion relation was only published by Szegő in 1939, so he didn't have a form in term's of  $\alpha_n$  and  $\rho_n$ .

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# Szegő's Theorem: Szegő-Verblunsky sum rule

In two remarkable 1935-36 papers, long unappreciated, Samuel Verblunsky (then just past his PhD. under Littlewood)

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$$\sum_{j=0}^{\infty} \log(1 - |\alpha_j|^2) = \int \log(w(\theta)) \frac{d\theta}{2\pi}$$

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It is critical that this always holds although both sides may be  $-\infty$ .

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It is critical that this always holds although both sides may be  $-\infty$ . This implies what I've called a "spectral theory gem"

$$\sum_{j=0}^{\infty} |\alpha_j|^2 < \infty \iff \int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$$

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In particular,  $\sum_{j=0}^{\infty} |\alpha_j|^2 < \infty \Rightarrow \Sigma_{ac} = \partial\mathbb{D}$ .

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# Szegő's Theorem: Szegő-Verblunsky sum rule

What makes the gems so interesting is that they allow arbitrary singular parts of the measures so long as the Szegő condition holds, i.e.  $\int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$ .

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What makes the gems so interesting is that they allow arbitrary singular parts of the measures so long as the Szegő condition holds, i.e.  $\int \log(w(\theta)) \frac{d\theta}{2\pi} > -\infty$ . If  $\sum_{j=0}^{\infty} |\alpha_j| < \infty$ , one can show that there is a scattering theory and strong asymptotic completeness holds in that there is only a.c. spectrum. The VS sum rules implies in going from  $\ell^1$  to  $\ell^2$  Verblunsky coefficients, one can have arbitrary mixed spectral types.

In the late 1990's unaware of the OPUC literature, my research group was studying  $1D$  Schrodinger operators,  $-\frac{d^2}{dx^2} + V(x)$  and the difference between  $L^1$  and  $L^2$  conditions.

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In the late 1990's unaware of the OPUC literature, my research group was studying  $1D$  Schrodinger operators,  $-\frac{d^2}{dx^2} + V(x)$  and the difference between  $L^1$  and  $L^2$  conditions. Deift-Killip had proven there was a.c. spectrum for  $L^2$  and showing there were examples with mixed spectrum was one of the problems in my list at the 2000 ICMP. Little did I know that an analogous problem had been solved in 1935!

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# Jacobi Parameters

Next orthogonal polynomials on the real line, aka OPRL. One starts with a probability measure,  $\mu$ , of compact support in  $\mathbb{R}$  and forms the orthonormal polynomials,  $\{p_n(x)\}_{n=0}^{\infty}$ . They obey recursion relations

$$xp_n(x) = a_{n+1}p_{n+1}(x) + b_{n+1}p_n(x) + a_np_{n-1}(x)$$

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$$xp_n(x) = a_{n+1}p_{n+1}(x) + b_{n+1}p_n(x) + a_np_{n-1}(x)$$

which sets up a one-one correspondence (which we'll call the *Jacobi map*) between such measures (with an infinity of points in their support) and sequences  $\{a_n, b_n\}_{n=1}^\infty$  of bounded  $a$ 's in  $(0, \infty)$  and  $b$ 's in  $\mathbb{R}$  (called *Jacobi parameters*).

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# Jacobi Parameters

There is also a correspondence between point measures with finite support and suitable sets of finitely many Jacobi parameters. If there are  $n$  pure points, then  $P_n$  is 0 in  $L^2(d\mu)$  so  $a_n = 0$  and again there are  $2n - 1$  Jacobi parameters –  $n$   $b$ 's and  $n - 1$   $a$ 's.

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For later purposes, I need some details on one approach to going from measures to Jacobi parameters. The more usual method than the one I want to discuss just forms the OPs and looks at the recursion parameters.

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For later purposes, I need some details on one approach to going from measures to Jacobi parameters. The more usual method than the one I want to discuss just forms the OPs and looks at the recursion parameters. Instead, consider the *once stripped Jacobi parameters*, i.e.  $\{a_{j+1}, b_{j+1}\}_{j=1}^{\infty}$  obtained by dropping the first row and column of the Jacobi matrix.

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For later purposes, I need some details on one approach to going from measures to Jacobi parameters. The more usual method than the one I want to discuss just forms the OPs and looks at the recursion parameters. Instead, consider the *once stripped Jacobi parameters*, i.e.  $\{a_{j+1}, b_{j+1}\}_{j=1}^{\infty}$  obtained by dropping the first row and column of the Jacobi matrix. For any non-trivial probability measure of compact support, let  $m(z) = \int d\mu(x)/(x - z)$  and let  $m_1$  be the spectral measure for the once stripped problem.

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# Jacobi Parameters

Then one can prove that

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# Jacobi Parameters

Then one can prove that

$$m(z) = \frac{1}{b_1 - z - a_1^2 m_1(z)}$$

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# Jacobi Parameters

Then one can prove that

$$m(z) = \frac{1}{b_1 - z - a_1^2 m_1(z)}$$

Using that  $m_1(z) = -z^{-1} + O(z^{-2})$ , one sees that one can go from the measure to  $m$

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Using that  $m_1(z) = -z^{-1} + O(z^{-2})$ , one sees that one can go from the measure to  $m$  to  $a_1, b_1$  and  $m_1$

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One consequence of this is that the poles of  $m_1$  (i.e. the pure points of  $\mu_1$ ) are precisely the zeros of  $m$ .

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# Szegő Condition

Here is one version of Szegő's Theorem for OPRL.

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# Szegő Condition

Here is one version of Szegő's Theorem for OPRL. The map  $z \mapsto z + z^{-1}$  maps  $\partial\mathbb{D}$  to  $[-2, 2]$  (via  $e^{i\theta} \mapsto 2 \cos \theta$ )

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$$\liminf_{n \rightarrow \infty} \prod_{j=1}^n a_j = \sqrt{2} \exp \left( \int_{-2}^2 \log |\pi s(x) w(x)| s(x) \frac{dx}{4\pi} \right)$$

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The condition for the finiteness of the integral is called the *Szegő condition*:

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The condition for the finiteness of the integral is called the *Szegő condition*:

$$\int_{-2}^2 \log |w(x)| (4 - x^2)^{-1/2} dx > -\infty$$

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# Szegő Condition

This doesn't yield a gem because

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This doesn't yield a gem because

$$\inf_n \prod_{j=1}^n a_j > -\infty \iff \int_{-2}^2 \log |w(x)| (4 - x^2)^{-1/2} dx$$

only holds under the a priori condition that  $\mu$  is supported inside  $[-2, 2]$

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only holds under the a priori condition that  $\mu$  is supported inside  $[-2, 2]$  and this is not simply expressible in terms of the Jacobi parameters; for example, it doesn't only depend on the parameters near  $\infty$  and can be changed by modifying a single  $a$  or  $b$ .

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# Killip–Simon Theorem

In 2001 (published 2003), Killip and I proved the following gem which we regard as an OPRL analog of the Verblunsky–Szegő gem where  $\{E_j^\pm\}_{j=1}^{N_\pm}$  are the eigenvalues outside  $[-2, 2]$  (with  $+$  above 2 and  $-$  below -2):

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**Killip–Simon Theorem** If  $d\mu = w(x)dx + d\mu_s$  is a measure of compact support on  $\mathbb{R}$  and  $\{a_n, b_n\}_{n=1}^\infty$  its Jacobi parameters, then

$$\sum_{j=1}^{\infty} |a_j - 1|^2 + b_j^2 < \infty$$

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# Killip–Simon Theorem

In 2001 (published 2003), Killip and I proved the following gem which we regard as an OPRL analog of the Verblunsky–Szegő gem where  $\{E_j^\pm\}_{j=1}^{N_\pm}$  are the eigenvalues outside  $[-2, 2]$  (with  $+$  above 2 and  $-$  below -2):

**Killip–Simon Theorem** If  $d\mu = w(x)dx + d\mu_s$  is a measure of compact support on  $\mathbb{R}$  and  $\{a_n, b_n\}_{n=1}^\infty$  its Jacobi parameters, then

$$\sum_{j=1}^{\infty} |a_j - 1|^2 + b_j^2 < \infty$$

if and only if the essential support of  $\mu$  is  $[-2, 2]$  and

$$\int_{-2}^2 \log(w(x)) \sqrt{4 - x^2} dx > -\infty \quad \sum_{j,\pm} (|E_j^\pm| - 2)^{3/2} < \infty$$

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# Killip–Simon Theorem

This result on Jacobi Hilbert-Schmidt perturbations of the free Jacobi matrix should be compared with a celebrated theorem of von-Neumann that any bounded self-adjoint operator has a Hilbert-Schmidt perturbation with only dense point spectrum!

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We called  $\int_{-2}^2 \log(w(x)) \sqrt{4-x^2} dx > -\infty$  *the quasi-Szegő condition* since the square root appeared to the  $+1/2$  power rather than the  $-1/2$  in the Szegő condition.

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$$\sum |E_n|^p \leq C \int |V(x)|^{p+d/2} dx$$

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# Killip–Simon Sum Rule

The gem comes from a sum rule. Let

$$Q(\mu) = \frac{1}{2\pi} \int_0^{2\pi} \log \left( \frac{\sin(\theta)}{\operatorname{Im} m(2 \cos(\theta))} \right) \sin^2(\theta) d\theta,$$

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$$F(E) \equiv \frac{1}{4}[\beta^2 - \beta^{-2} - \log(\beta^4)] \quad E = \beta + \beta^{-1} \quad |\beta| > 1$$

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Then the Killip–Simon sum rule says

$$Q(\mu) + \sum_{j,\pm} F(E_j^\pm) = \sum_{n=1}^{\infty} \frac{1}{4} b_n^2 + \frac{1}{2} G(a_n)$$

As with the Szegő–Verblunsky sum rule, an important point is that it always holds although both sides may be  $+\infty$ .

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The gem comes from the fact that  $F \geq 0$ , vanishes exactly at  $E = \pm 2$  and is  $O((|E| - 2)^{3/2})$  there and that  $G \geq 0$ , vanishes exactly at  $a = 1$  and is  $O((a - 1)^2)$  there.

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The positivity of the terms is essential to be sure that there aren't cancelations.

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As in the OPUC case, this sum rule implies the existence of Hilbert-Schmidt perturbations with mixed spectrum.

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# Step-by-Step Sum Rule

I know of many proofs of Szegő's Theorem but until recently all proofs of the Killip–Simon sum rule were variants of our original proof which I want to describe some parts of.

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Because the eigenvalues  $E_j^\pm$  and  $E_j^{(1)\pm}$  interlace and  $F$  is monotone, the sum is of positive terms and always convergent (interlacing sums).

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While the  $\log(w)$  integral might diverge, one can show that a  $\log(w/w_1)$  integral is always convergent.

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If you iterate coefficient striping and assume the boundary term goes away, you get the full sum rule.

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# Poisson–Jensen Formula

The step-by-step sum rule will involve a Poisson–Jensen formula whose classical form we recall.

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The step-by-step sum rule will involve a Poisson–Jensen formula whose classical form we recall. Define Blaschke factors,  $b(z, w)$  to be  $z$  is  $w = 0$  and otherwise  $-\frac{|w|(z-w)}{w(1-\bar{w}z)}$ . Let  $f$  be analytic on the unit disk and in Nevanlinna class, i.e.  $\sup_{0 < r < 1} \int \log_+(|f(re^{i\theta})|) d\theta < \infty$ . If  $\{z_j\}_{j=1}^N$  is a listing of the zeros of  $f$ , then  $\sum_{j=1}^N (1 - |z_j|) < \infty$  which implies that  $B(z) = \prod_{j=1}^N b(z, z_j)$  converges to an analytic function vanishing exactly at the  $z_j$ . Suppose also that for some  $p > 1$ , we have that  $\log(f(z)/B(z))$  lies in  $H^p$  (which we'll call the " $L^p$ -condition").

The famous theorem of Smirnov and Beurling says that for some  $\omega \in \partial\mathbb{D}$ , we have that (Poisson–Jensen formula)

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Without the  $L^p$ -condition, there a singular inner part.





# Meromorphic Herglotz Functions

By a *meromorphic Herglotz* function, we mean a function meromorphic on  $\mathbb{D}$ , real on  $(-1, 1)$  with  $\operatorname{Im} z > 0 \Rightarrow \operatorname{Im} f(z) > 0$ .

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# Meromorphic Herglotz Functions

One can prove that in  $\mathbb{D} \cap \mathbb{C}_+$ , one has that  $|\arg z B(z)| \leq 2\pi$  so that  $\arg(f(z)/zB(z))$  is bounded on  $\mathbb{D}$ .

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Taking log's, one gets relations between Taylor coefficients of  $\log(f(z)/z)$ , certain sums involving logs or powers of zeros and poles and integrals  $\cos(n\theta) \log |f(e^{i\theta})|$ .

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# Case Step-by-Step Sum Rules

Recall that  $m(z) = \int d\mu(x)/(x - z)$ . It defines a Herglotz function on  $\mathbb{C}_+$ , real on  $\mathbb{R}$ .

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The above procedure thus yields a relation between polynomials of Jacobi parameters, the difference of functions of the eigenvalues of  $J$  and  $J_1$  and integral of  $\log |M(e^{i\theta})|$ .

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# $P_2$ Sum Rule

What results is a step-by-step sum rule which if iterated with boundary terms dropped yields the formal sum rules stated by Case

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## $P_2$ Sum Rule

What results is a step-by-step sum rule which if iterated with boundary terms dropped yields the formal sum rules stated by Case (although, unlike Case, Killip and I had explicit formulae for the polynomials in the Jacobi parameters).

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However, we discovered that  $C_0 + \frac{1}{2}C_2$  had the required positivity. We had no explanation of why this was so but observed it. We called this the  $P_2$  sum rule (P for positive) and it is now known as the Killip–Simon sum rule.

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However, we discovered that  $C_0 + \frac{1}{2}C_2$  had the required positivity. We had no explanation of why this was so but observed it. We called this the  $P_2$  sum rule (P for positive) and it is now known as the Killip–Simon sum rule. The rather complicated functions  $F$  and  $G$  just arose by taking the functions from the Case sum rule and combining them.

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# Mysteries

While the gem one gets from the  $P_2$  sum rule is simple and elegant, the proof has lots of mysteries:

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While the gem one gets from the  $P_2$  sum rule is simple and elegant, the proof has lots of mysteries:

- ① Why are there any positive combinations?
- ② It is easy to understand the  $(4 - x^2)^{-1/2} dx$  of the Szegő condition. It is  $d\theta$  under  $x = \cos(\theta)$ . Equivalently, it is the potential theoretic equilibrium measure for  $[-2, 2]$

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- ③ What does the function

$$G(a) = a^2 - 1 - \log(a^2)$$

mean?

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# Mysteries

While the gem one gets from the  $P_2$  sum rule is simple and elegant, the proof has lots of mysteries:

- ① Why are there any positive combinations?
- ② It is easy to understand the  $(4 - x^2)^{-1/2} dx$  of the Szegő condition. It is  $d\theta$  under  $x = \cos(\theta)$ . Equivalently, it is the potential theoretic equilibrium measure for  $[-2, 2]$  but where the heck does the  $(4 - x^2)^{1/2} dx$  come from?

- ③ What does the function

$$G(a) = a^2 - 1 - \log(a^2)$$

mean?

- ④ What does the function

$$F(E) = \frac{1}{4}[\beta^2 - \beta^{-2} - \log \beta^4]; \quad E = \beta + \beta^{-1}$$

mean?

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# The LD Framework

Large deviations go back to Laplace. The modern theory was initiated by Cramér in the 1930's and made into a powerful machine by Donsker–Varadhan and Freidlin–Wentzel and then Varadhan alone (work for which he got the Abel prize).

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We consider a sequence of probability measures,  $\{\mu_n\}_{n=1}^\infty$ , on a space,  $X$ . Naively, one has a Large Deviation Principle (aka LDP) if the  $\mu_n$ -probability that  $x$  is near  $x_0$  is  $O(e^{-nI(x_0)})$ .

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# The LD Framework

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# The LD Framework

①  $I$  is lower semicontinuous

② For all closed sets  $F \subset X$

$$\limsup_{n \rightarrow \infty} \frac{1}{a_n} \log \mu_n(F) \leq -\inf_{x \in F} I(x)$$

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$$I(x) = \sup_{\theta} \left[ \theta x - \log \left( \mathbb{E}(e^{\theta X}) \right) \right]$$

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# Cramér Example

Let  $X$  be an exponential random variable, i.e. with density  $\chi_{[0,\infty)}(x)e^{-x} dx$ . Then

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$$I(x) = \begin{cases} x - 1 - \log(x), & \text{if } x > 0 \\ \infty, & \text{if } x \leq 0 \end{cases}$$

Notice that  $G(a) = I(a^2)$ ,

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# LDP and Sum Rules

Gamboa, Nagel and Rouault had the following lovely idea.  
Let  $X$  be the set of probability measures on  $\partial\mathbb{D}$  or on  $\mathbb{R}$   
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# LDP and Sum Rules

GNR had the further idea that the measures on the spectral measures should come from random matrix measures with a cyclic vector in the limit as the matrix dimension goes to infinity.

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GNR had the further idea that the measures on the spectral measures should come from random matrix measures with a cyclic vector in the limit as the matrix dimension goes to infinity.

Of course, the issue becomes to effectively compute the rate function on both sides and alas, we haven't yet found a magic way to do these calculations in a general context.

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The reception of the GNR paper illustrates the dangers of working in between two disparate areas.

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The reception of the GNR paper illustrates the dangers of working in between two disparate areas. They wrote the paper in a way that only experts on large deviations could understand it, but such experts didn't understand the spectral theory context.

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# LDP and Sum Rules

Jonathan Breuer and I couldn't understand the paper so we consulted Ofer Zeitouni who said he'd looked quickly at the paper and there didn't seem to be much new there! In fact, the calculations of rate functions on the two sides wasn't so far from prior calculations of rate functions. What was new was the realization that because a rate function could be computed in two ways, one is able to prove interesting equalities. So they had some troubles getting published what I regard as one of the more interesting recent papers in spectral theory. In the end, Jonathan, Ofer and I used their methods to study higher order sum rules

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# CUE: Measure Side

Circular Unitary Ensemble, aka CUE, is just another name for Haar Measure in  $U(n)$ , the  $n \times n$  unitary matrices, for varying  $n$ .

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As is well-known, the  $\lambda$ 's and  $w$ 's are independent of each other, the  $w$ 's are uniformly distributed on the simplex  $\{\mathbf{w} \mid \sum_{j=1}^n w_j = 1\}$

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As is well-known, the  $\lambda$ 's and  $w$ 's are independent of each other, the  $w$ 's are uniformly distributed on the simplex  $\{\mathbf{w} | \sum_{j=1}^n w_j = 1\}$  and by the Weyl integration formula, the  $\theta$ 's have distribution

$$\frac{1}{n!} \prod_{1 \leq j < k \leq n} |e^{i\theta_j} - e^{i\theta_k}|^2 \prod_{j=1}^n \frac{d\theta_j}{2\pi}$$

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# CUE: Measure Side

The first step in the analysis of the measure side is to analyze what probabilists call the *empirical measure* and physicists *the density of states*, namely the random measure

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This is easy to understand. The Weyl distribution can be viewed as a discrete two dimensional Coulomb gas in the canonical ensemble

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This is easy to understand. The Weyl distribution can be viewed as a discrete two dimensional Coulomb gas in the canonical ensemble (2D because  $|x - y|^{-2}$  is the exponential of  $-2 \log |x - y|$ ). The  $n \rightarrow \infty$  limit is a high density limit and due to repulsion, there is a strong tendency towards equal spacing.

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# CUE: Measure Side

To get a significant difference from equal spacing, one has  $O(n^2)$  smaller distances and so the speed is  $n^2$ . The optimal spacing will still be locally equal and the discrete Coulomb energy will converge to the continuum.

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The fact that  $n^2$  is much larger than  $n$  implies that for a measure to have finite rate at speed  $n$ , it has to have points close to uniformly distributed

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The weights are close to independent (except for the normalization they are)

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The weights are close to independent (except for the normalization they are) – a slick way to see this is to note if  $Y_j$  are positive exponentially distributed iidrv, then  $w_j = Y_j / \sum_{j=1}^n Y_j$ . This allows one (using Cramér's theorem on small blocks) to prove an LDP for the spectral measure with speed  $n$  and rate function the Szegő integral  $-\int \log(w(\theta)) \frac{d\theta}{2\pi}$ .

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# CUE: Coefficient Side

In 2004, Killip and Nenciu wrote down the distribution of  $\{\alpha_j\}_{j=0}^{n-1}$  induced by restricting Haar measure as we are.

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$$\frac{n-j-1}{\pi} (1 - |z|^2)^{n-j-2} d^2z$$

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which says that  $\alpha_j$  is distributed as the first complex component of a unit vector in  $\mathbb{C}^{n-j}$ .

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# GUE

The Gaussian Unitary Ensemble, aka GUE, is the probability measure on  $n \times n$  self adjoint matrices so that  $\{a_{ii}\}_{i=1,\dots,n}$ ,  $\{\operatorname{Re}(a_{ij})\}_{1 \leq i < j \leq n}$  and  $\{\operatorname{Im}(a_{ij})\}_{1 \leq i < j \leq n}$  are independent identically distributed Gaussian random variables of mean zero and suitable,  $n$ -dependent variance.

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The argument for GUE, normalized so the limiting density is the semicircle law on  $[-2, 2]$ , is similar to that for CUE.

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The argument for GUE, normalized so the limiting density is the semicircle law on  $[-2, 2]$ , is similar to that for CUE. Instead of results of Killip-Nenciu for the distribution of  $\alpha$ 's, one has earlier results of Dumitriu and Edelman for the Jacobi parameters. The calculation is made easier by the independence of the Jacobi parameters (which leads to sums of terms that depend only on a single  $a$  or  $b$ ).

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# GUE

One needs to make some additional arguments going back to Ben Arous-Dembo-Guionnet to deal with eigenvalues outside the essential support.

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# GUE

One needs to make some additional arguments going back to Ben Arous-Dembo-Guionnet to deal with eigenvalues outside the essential support.

What results is a new proof of the Killip-Simon sum rule.

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*mean?*

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*mean?* This is the Coulomb potential of the Wigner semi-circle distribution plus a quadratic external field.

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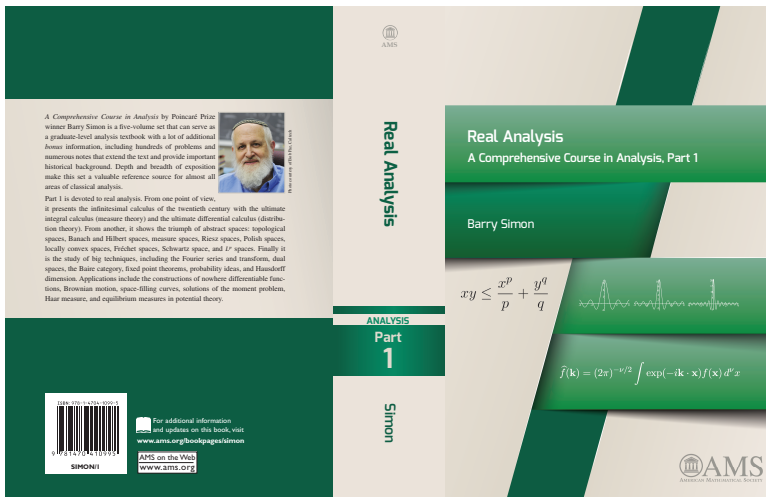
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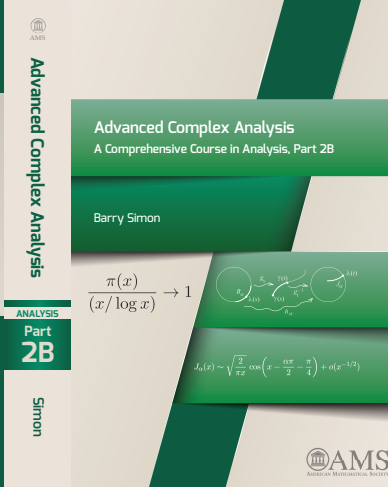



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**Part 3**

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**Harmonic Analysis**

**A Comprehensive Course in Analysis, Part 3**

**Barry Simon**



$$|f - f_Q|_Q = \frac{1}{|Q|} \int_Q |f(x) - f_Q| dx$$

$$|(x \mid M_{H_L} f(x) > \alpha)| \leq \frac{3^{\alpha}}{\alpha} \|f\|_{L^1(\mathbb{R}^n, d^x)}$$

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*Part 3 returns to the themes of Part 1 by discussing pointwise limits (going beyond the usual focus on the Hardy–Littlewood maximal function by including ergodic theorems and martingale convergence), harmonic functions and potential theory, frames and wavelets,  $H^p$  spaces (including bounded mean oscillation (BMO)) and, in the final chapter, lots of inequalities, including Sobolev spaces, Calderón–Zygmund estimates, and hypercontractive semigroups.*



Photograph of Barry Simon



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Part 4 focuses on operator theory, especially on a Hilbert space. Central topics are the spectral theorem, the theory of trace class and Fredholm determinants, and the study of unbounded self-adjoint operators. There is also an introduction to the theory of orthogonal polynomials and a long chapter on Banach algebras, including the commutative and non-commutative Gelfand–Naimark theorems and Fourier analysis on general locally compact abelian groups.

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Operator Theory

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$$A = \int t dE_t$$

$$\det(1 + zA) = \prod_{k=1}^{N(A)} (1 + z\lambda_k(A))$$

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